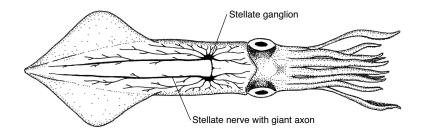
Response of the Hodgkin-Huxley model in the high-input regime

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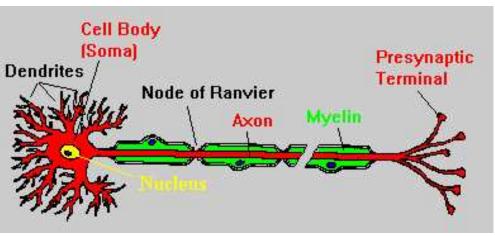


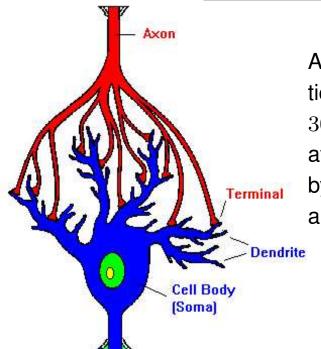
Summary

- Brief introduction of the HH model
- Characterization of the stochastic stimulation protocol
- Analysis of the neuronal responses for different noise levels
- Looking for coherence in the neuronal response
- Influence of correlations on the coherent response
- Conclusions and open problems



The neuron in brief





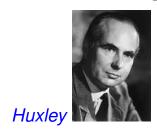
A neuron in the brain cortex has many ($\sim 10,000$) synaptic connections, but not all active. The neuron receives 300-1,000 postsynaptic inputs of amplitude $\sim 0.5-1$ mV at a frequency ~ 100 Hz, it responds each 10-40 inputs by emitting an action potential of duration 1-2 msec and amplitude ~ 100 mV.





Hodgkin

The HH model reproduces the time evolution of the membrane potential and of the ionic currents measured experimentally for a giant squid axon.

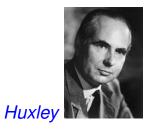




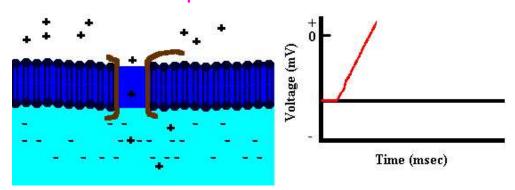


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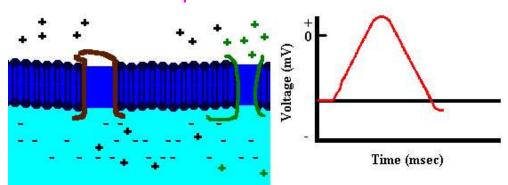


Membrane Depolarization



 Na^+ enters inside the cell

Membrane Repolarization



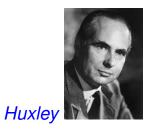
 K^+ leaves the cell

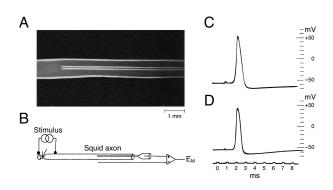




Hodgkin

The HH model reproduces the time evolution of the membrane potential and of the ionic currents measured experimentally for a giant squid axon.





 $C=1\mu F/cm^2$ - Membrane capacitance

V - Membrane Potential (mV)

 I_i - Ionic channel currents $(\mu A/cm^2)$

$$C\dot{V} = \sum_{i} I_{j} + I_{syn} = -g_{Na}m^{3}h(V - V_{Na}) - g_{K}n^{4}(V - V_{K}) - g_{L}(V - V_{L}) + I_{syn}$$

$$\dot{x} = \alpha_x - x(\alpha_x + \beta_x)$$
 $x = n, m, h$ gating variables

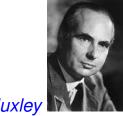
 $\alpha_x = \alpha_x(V)$ and $\beta_x = \beta_x(V)$ are highly nonlinear functions.





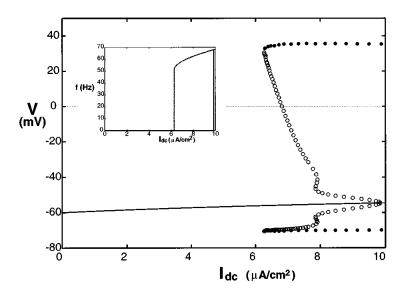
Hodgkin

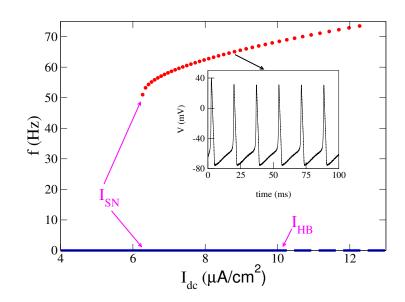
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Huxley

Constant Current Synaptic Input $I_{syn} = I_{dc}$







Instead of a constant current I_{dc} , we consider N_E excitatory (EPSP) and N_I inhibitory postsynaptic inputs (IPSP), each corresponding to a voltage kick $\Delta V = 0.5$ mV.



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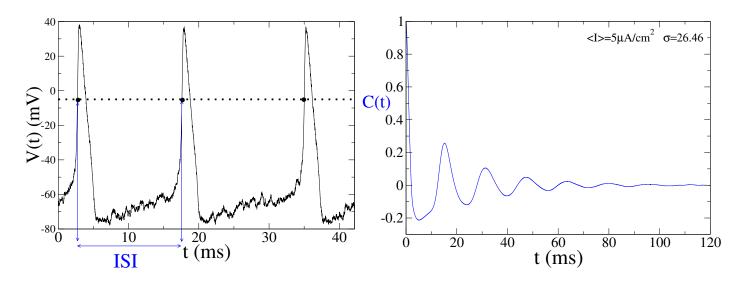
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At these frequencies the net input spike count within a temporal window ΔT (≥ 1 msec) is essentially Gaussian distributed and it can be characterized by its average $\mu = \nu (N_E - N_I) \Delta T$ and variance $V = \nu (N_E + N_I) \Delta T = \nu \sigma^2 \Delta T$.

The response of the neuron is examined for fixed average input current $\bar{I} = C\Delta V \nu (N_E - N_I)$ by varying only σ and therefore the standard deviation of the noise.



Statistical and dynamical indicators

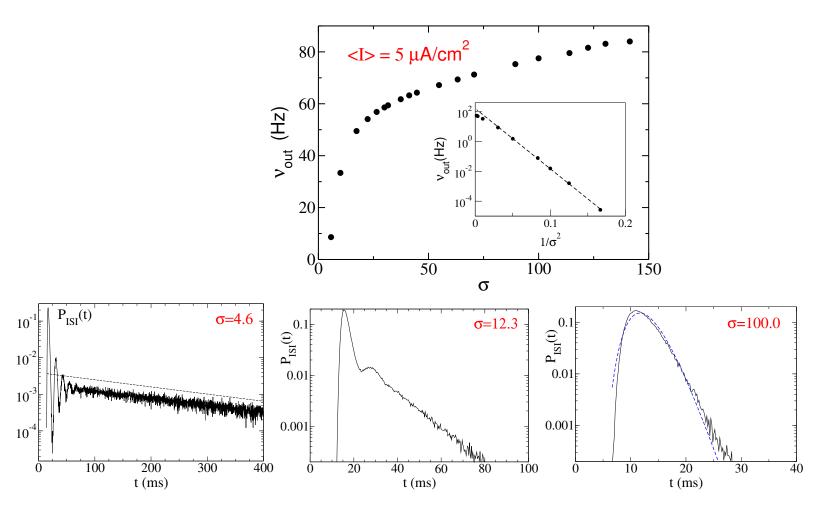


- **■** ISI distribution $\rightarrow P_{ISI}(t)$;

- Conditional entropies (not discussed here)

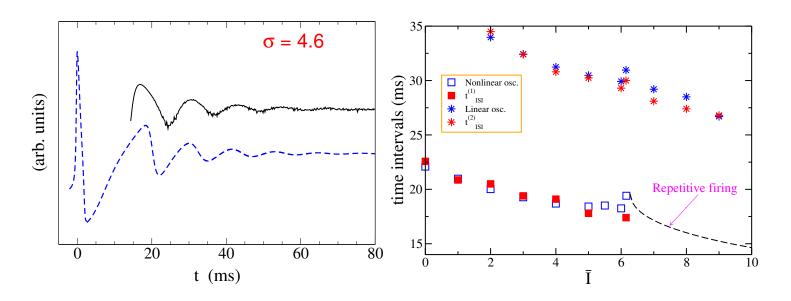


The HH neuron is in the silent state, i.e. the average input current \bar{I} is smaller than I_{SN} .





Spikes triggered by relaxation oscillations



The neuronal spiking is enhanced in correspondence of the maxima of the relaxation oscillations following a spike emission.

The first oscillation has a nonlinear origin, while the period of the subsequent ones can be obtained via a linear analysis around the stable fixed point solution (focus).



Firing activated by noise

Two mechanisms compete:

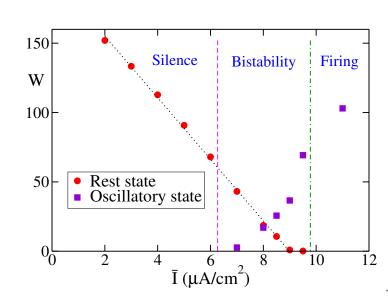
- the HH dynamics tends to relax towards the rest state;
- noise fluctuations lead the system towards an excitation threshold.

The dynamics of V(t) resembles the overdamped dynamics of a particle in a potential well under the influence of thermal fluctuations, and the firing times can be expressed in terms of the Kramers expression (for sufficiently small noise)

$$t_a \propto \mathrm{e}^{W_S/\sigma^2}$$

the time distribution is Poissonian (CV = 1).

- for $\sigma < \sqrt{W_S} \rightarrow$ Activation Process
- for $\sigma > \sqrt{W_S} \rightarrow \text{Diffusive Dynamics}$





High noise limit

The effect of noise fluctuations on the neuron dynamics is twofold:

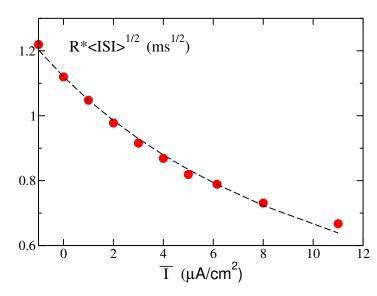
- **\blacksquare** a constant current \bar{I} driving the system;
- a stochastic term with zero average.

The dynamics of V(t) can therefore be described in terms of a Langevin process with a drift and the distribution of the first passage times is given by the inverse Gaussian distribution:

$$f(t) = \frac{\alpha}{\sqrt{2\pi\beta t^3}} e^{-\frac{(t-\alpha)^2}{2\beta t}}$$

In this case the coefficient of variation should be given by

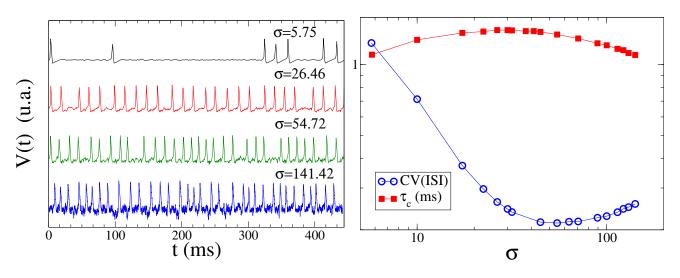
$$CV \propto \frac{\sigma}{(\bar{I} + I_0)\sqrt{\langle ISI \rangle}}$$





Coherence resonance

Coherence of the emitted spike trains



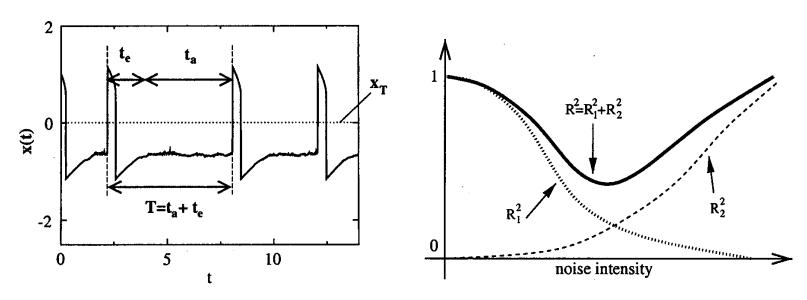
In the silent and bistable regime ($I < 8\mu\text{A/cm}^2$):

- \blacksquare in the beginning the HH neuron response becomes more regular for increasing σ ;
- the maximal coherence is reached for a finite σ -value;
- for higher noise amplitudes the response becomes again more irregular.

A. Pikovsky & J. Kurths, PRL 78, 775 (1997)



Coherence resonance



- ullet The system is characterized by two characteristic times $ightarrow ISI \equiv T = t_a + t_e$:
 - **•** t_a =activation time \rightarrow time needed to excite the system;
 - t_e =excursion time \rightarrow duration of the spike (excited state).
- The competition of the two effects leads to an intermediate regime of coherence:

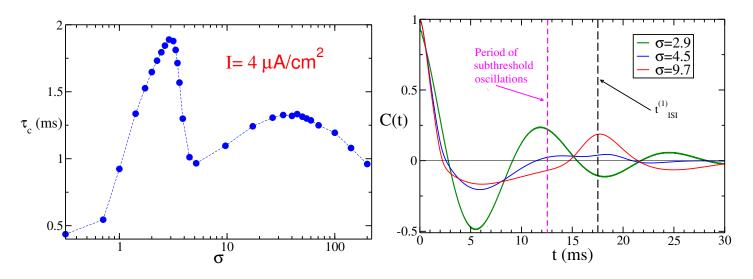
$$CV(ISI)^2 = CV(t_a)^2 \frac{\langle t_a \rangle^2}{\langle T \rangle^2} + CV(t_e)^2 \frac{\langle t_e \rangle^2}{\langle T \rangle^2} = R_1^2(t_a) + R_2^2(t_e)$$

 $R_1^2(t_a)$ decreases with σ , while $R_2^2(t_e)$ increases \rightarrow minimum in R(ISI)



Coherence resonance

Coherence of the subthreshold oscillations



A second coherence resonance is revealed by analyzing the correlations of the potential:

- for σ < 3 almost no spikes are emitted, but the increase of noise leads to more and more regular subthreshold oscillations;
- for $\sigma > 3$ the statistics of the emitted spikes is no more negligible and this decorrelates the signal;
- for $\sigma > 10$ the dynamics is dominated by sequences of spikes and a second peak occurs related to the regularization of the spike trains.



Correlations generated via common drive

In order to generate N_E correlated (ρ) Poissonian spike trains (with rate ν_0):

- firstly a long sequence of M ISI's is extracted from a Poissonian distribution with rate ν_0/ρ ;
- secondly from this common pool $N_{sp} = M\rho$ ISI's are randomly extracted N_E times;
- the superposition of the N_E correlated trains gives rise to a sequence of kicks of variable amplitude (with average $\rho N_E \Delta V$) and with ISIs Poissonian distributed with rate ν_0/ρ .

The correlation between two input spike trains originating from neuron i and j is measured in terms of the Pearson correlation coefficient :

$$\rho = \frac{\langle (n_i - \langle n_i \rangle)(n_j - \langle n_j \rangle) \rangle}{s^2}$$

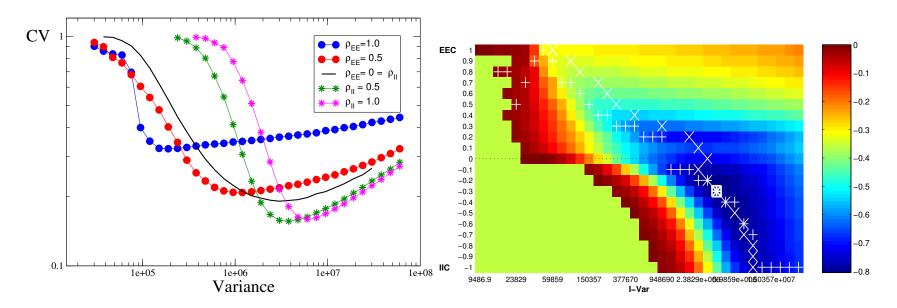
where n is the number of spikes in a time window ΔT and s^2 its variance.

M.N. Shadlen & W.T. Newsome (1998) – E. Salinas & J. Sejnowski (2000)



We have studied the response of the (balanced $N_E=N_I$) Fitz-Hugh Nagumo model in the silent regime for excitatory (resp. inhibitory) correlated inputs at constant correlation ρ by varying the noise variance. The increase of the variance leads to an increase of the average amplitude for the correlated kicks, while their frequency remains constant.

- Coherence resonance is observed for any excitatory (resp. inhibitory) correlation at finite noise amplitude.
- An absolute resonance with respect to noise and correlation can be identified.





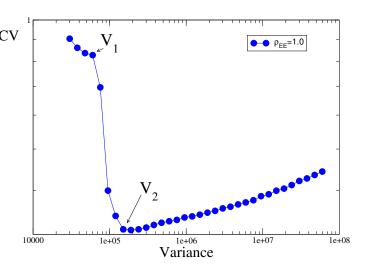
Strong excitatory correlation

The coherence phenomenon is now determined only by the kick amplitude and not by the properties of the asymptotic stochastic processes, since in the present case the output can be always described as a Poissonian process with a refractory time.

For increasing variance the amplitude of the correlated kicks increases.

- for $Variance < V_1 \rightarrow Activation Process$
- for $V_1 < Variance < V_2$ One kick is sufficient to induce a spike emission
- for $V_2 < Variance$ One kick elicits always a spike T_{ref} decreases

$$CV = \frac{T}{T + T_{ref}}$$



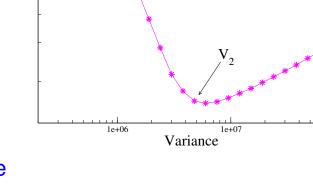
For $Variance > V_2 \rightarrow T = 1/\nu_0$, we have a 1:1 correspondence between input and output (apart from the refractory period).



Strong inhibitory correlation

At large variances the dynamics is again ruled by the the kick amplitude of the correlated kicks, but at lowervariances the correlated kicks are quite rare and their amplitude is not sufficient to influence the dynamics.

- for $Variance < V_1$ Activation Process $CV \simeq 1$
- for $V_1 < Variance < V_2$ Dynamics dominated by uncorrelated excitatory kicks leading the system in the repetitive firing regime
- for $Variance > V_2$ Each kick induces a certain delay in the spike time of the neuron - A multimodal structure appears in the ISI distribution



Frequency of the correlated kicks (ν_0) << Frequency of the uncorrelated kicks $(N\nu_0)$



Summary and perspectives

Uncorrelated stochastic inputs

- The response of a HH neuron depends only on the average and the variance of the input, at least in the high input regime;
- at $I < I_{SN}$ the neuronal firing, induced by the stochastic inputs, can be expressed as an activation process at low variances (σ^2), while for large σ^2 this process becomes essentially diffusive;
- at low noise, beside of the exponential tail, the ISI distributions reveal a multimodal structure due to spiking triggered by relaxation oscillations towards the rest state;
- coherence resonance can be observed in a large interval of currents in the silent and bistable regime;
- a second coherence resonance (associated to subthreshold oscillations) coexists with the usual one;

Correlated stochastic inputs

- new mechanisms for the coherence resonance have been reported at high excitatory and inhibitory correlations;
- coherence effects can be induced by varying only the correlation.



Summary and perspectives

What to do next?

Experiments

- Is it possible to reveal both kinds of CR with dynamic clamp measurements of Stellate Cells?
- Due to the high degree of erraticity in the response of the Stellate Cells, maybe CR could be observed without injecting noise in the sytem, but just varying the DC current.

Networks

- ▲ A single HH neuron can fire at different frequencies (when stochastically stimulated) therefore a network of coupled HH neurons should exhibit coherent and correlated activities over different time scales. (several rithms)
- How CR will influence network dynamics?



Credits

- Stefano Luccioli Msc in Physics (2004-2005)
- Dynamics of realistic single neuronal models

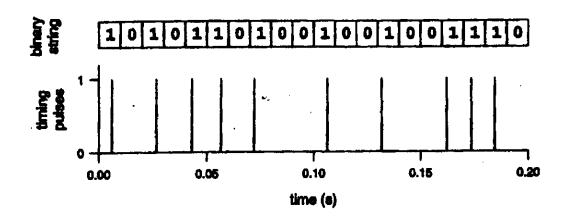


- Thomas Kreuz Marie Curie Fellow (2005-2006)
- Dynamical Entropies in Assemblies of Neurons





Entropie condizionali



- $\triangle t$ ="finestra" temporale \rightarrow *codifica binaria* ("1"/"0") del potenziale di membrana;
- $m{P}$ $C_N = (1,0,1,\ldots) \rightarrow \text{"parola"}$ (o "stato") di lunghezza N;
- lacksquare $H(N) = -\sum_{\{C_N\}} P(C_N) \log_2 P(C_N)$, "entropia del blocco N";
- $m{P}$ h(N) = H(N+1) H(N), "entropia condizionale" \rightarrow regolarità, prevedibilità:

$$(1,0,0,1,0,1,?,...)$$
 $h(N+1) \le h(N)$

$$h_{max}(N) = \log_2 2 = 1.$$

